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ANALYTICAL CORE LOSS CALCULATIONS
FOR MAGNETIC MATERIALS USED IN
HIGH FREQUENCY HIGH POWER
CONVERTER APPLICATIONS

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by James E. Triner

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Dean of the Graduate School

The University of Toledo August 1979

An Abstract of

ANALYTICAL CORE LOSS CALCULATIONS FOR MAGNETIC MATERIALS USED

IN HIGH FREQUENCY HIGH POWER CONVERTER APPLICATIONS

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Analysis of future space energy needs has indicated a variety of missions requiring total energies exceeding those of past missions by two or three orders of magnitude (1 to 100 kW). A major part of such a system will be the power converter and inverter subsystems which process this energy to meet user requirements.

One of the major components within such a power subsystem is the magnetic device. This component is used for voltage and current transformation, input/output filter sections, parameter sensing, and semiconductor control. Recently, attention has been given to the different loss characteristics of magnetic devices observed under various excitation conditions.

This dissertation examines the basic magnetic properties under various operating conditions encountered in the state-of-the-art DC-AC/DC converters. Using a novel core excitation circuit, the basic B-H and loss characteristics of various core materials may be observed as a function of circuit configuration, frequency of operation, input voltage, and pulse-width modulation conditions. From this empirical data, a mathematical loss characteristics equation is developed to analytically predict the specific core loss of several magnetic materials under various waveform excitation conditions.

Previous magnetic core characteristics were observed using either a voltage source or a current source. For a core being excited in either of these modes, the voltage or current characteristics are observed as the independent variable. However, the excitation circuit used for this dissertation uses a high power wide bandwidth operational amplifier which senses the derivative of flux with time as the controlled parameter. Under these conditions flux is observed as the independent variable with voltage and current as the dependent variable.

Power measurements were obtained using a true RMS wide bandwidth wattmeter. The B-H loop characteristics were observed using an X-Y oscilloscope. This experimental excitation set-up allowed the magnetic characteristics to be measured and observed under square, ramp, and sinusoidal flux conditions.

Using this flux controlled excitation, magnetic core characteristics were developed for constant values of ramp flux and frequency. These characteristics will show the circuit designer for the first time the functional relationships between induction level and specific core loss as a function of the two key DC-DC converter operating parameters of input voltage and duty cycle.

Core excitations involving symmetric and asymmetric conditions for both sinusoidal and square wave waveforms were investigated analytically using three specific core loss calculation methods. Experimental data was used to compare the accuracy of these methods. The predicted and measured specific core losses were within 10 percent of each other for all conditions analyzed. This deviation is within the error band of the instrumentation. Therefore, these methods will analytically predict the specific core loss for any type of voltage excitation encountered in DC-AC/DC converter configurations using only one set of characteristic curves for the specific magnetic material, core configuration, and tape thickness.

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CHAPTER I

INTRODUCTION

Magnetic materials offer the power electronics designer a widerange of magnetic characteristics (Ref. 1) from which he must select those materials applicable to his particular DC-AC/DC converter design. Like their semiconductor counterparts, magnetic devices display families of operating characteristics which are functions of frequency, temperature, basic material composition, fabrication techniques, and driving source impedance (Ref. 2). The interaction between these various functions may be explained by ferromagnetism, material properties, or domain wall movement (Ref. 3). However, for power electronics applications it is desirable to provide a set of magnetic characteristics which are easily understood and accurately predict the performance of a particular magnetic material when applied in a specific DC-AC/DC converter configuration.

This dissertation examines the B-H and loss characteristics of several magnetic materials used in present state-of-the-art power processing circuits. The effects of variables such as frequency, voltage level, voltage waveforms, and modulation upon the inherent material characteristics are explored using only externally measurable quantities.

Traditionally derived core loss data are presented as a function of flux density with frequency as a parameter. In general, hard experimental data is taken with sinusoidal voltage excitation of 60 or

erated by using a mathematical formula such as that proposed by Legg (Ref. 4). However, this method is only valid for low values of flux density. Only recently has attention been given to the different loss characteristics of magnetic materials observed between sinusoidal and square-wave excitation conditions at 10,000 hertz (Ref. 5). In the past, only circuits having similar excitation conditions could utilize these loss characteristics. Newly available state-of-the-int institutementation equipment, and mathematical techniques developed in this dissertation allow the power electronics designer to calculate the specific core loss for various nonsinusoidal, and asymmetrical square wave excitation conditions using only these loss characteristics.

The square wave B-H loop and associated core loss characteristics may be observed in a unique manner by means of an excitation source in which the flux rate of change $(d\phi/dt)$ is observed as the independent variable, and the power loss is measured using a wide bandwidth (DC-300 K Hz) wattmeter. This excitation circuit, which is fully described in Chapter II, allows the control of an additional variable, $d\phi/dt$, rather than just the driving frequency as is the case for all previous sinusoidal core loss measurements. If the excitation is restricted to that of a constant frequency square wave voltage having a variable duty-cycle, the $d\phi/dt$ effects may then be determined for various constant values of $B_{\rm M}$ independent of frequency. For the first time, the distribution of the B-H loop area is shown to be dependent upon the positive and negative $d\phi/dt$ ramp rates. The wideband wattmeter allows core loss data to be accurately measured under these conditions. This procedure allows a true separation of the variables d^*/dt , frequency,

and B_M . This data is then used to develop the functional relation-ships of nysteresis loss with B_M , and eddy current loss with B_M and $d\phi/dt$.

Sinusoidal core loss characteristics for a particular magnetic material and lamination thickness are generally shown with flux density as the independent variable, specific core loss as the dependent variable, and frequency as a parameter (Ref. 6). Data collected with the new excitation circuit is based upon values of constant $d\phi/dt$ which corresponds with square wave excitation. The core loss characteristics are snown with frequency as the independent variable, specific core loss as the dependent variable, and $d\phi/dt$ and flux density as parameters.

In Chapter III these loss characteristics, the B-H loop behavior, and the instantaneous power waveshapes observed for the various waveforms are used to develop three methods which analytically predict the specific core loss for magnetic materials. A pulse width modulated square wave with a 10 to 50 percent duty cycle is then used to compare the experiment. and analytical accuracy of each method.

The first method, which is similar to the sinusoidal analysis techniques already in common use, allows the specific core loss to be calculated directly from the characteristics knowing the maximum flux density $(B_{\underline{M}})$ and the frequency of operation. To calculate the value of $B_{\underline{M}}$ for the excitation pulse, the cross sectional area of the core, number of turns, excitation voltage level, and pulse width must be known. In the second method, the specific core loss is calculated by using $d\phi/dt$ and an "equivalent" excitation frequency.

The value of do/dt is directly proportional to the input voltage

and inversely proportional to the number of turns on the core. The "equivalent" frequency is inversely proportional to twice the pulse width. The third method is developed as an extension of the second method. The excitation pulse is subdivided into a discrete number of time steps. The average voltage level is determined for each time step. This voltage level along with the number of turns on the core correspond to an incremental value of d\$\phi/dt\$. This corresponds to an incremental loss at some "equivalent" frequency. The total specific core loss is then the summation of these incremental losses divided by total period. These methods cover the functional relationship between induction level and specific core loss as a function of excitation voltage level and duty-cycle characteristics.

The concepts developed for the square wave in the third method are extended to analyze other types of waveforms in Chapter IV. A general equation form is used to relate the specific core loss to d\$\psi/dt\$, and the equivalent frequency of the excitation. The coefficient and expenents that relate these variables are dependent on the magnetic material, and the tape thickness, or grain size. Using this general equation a computer algorithm is then developed which automates the analysis procedure. Analytical and experimental results are compared for sinusoidal and nonsinusoidal examples using three different core materials and configurations. As a final example the specific core loss characteristics for 1/2 mil Supermalloy material under sinusoidal excitation are generated using the characteristic equation for square wave excitation and the computer algorithm.

Finally, Chapter V presents the conclusions of this work and suggestions for future research in this area. Spec fic core loss data for several magnetic tape materials and ferrite are presented in Appendix A. A glossary of symbols and their associated units are in Appendix B.

CHAPTER II

INVESTIGATION OF MAGNETIC CORE CHARACTERISTICS

2.1 Introduction

The analytical core loss calculation methods developed in this dissertation are based upon the core characteristics observed under controlled flux excitation. The width of the B-H loop was found to be a function of both the flux rate of change $(d\phi/dt)$ and the maximum flux density (B_M) .

In all the specific core loss (SCL) characteristics presented, both positive and negative values of d\(\phi/\)dt were held constant over one cycle. These d\(\phi/\)dt conditions allow symmetrical square wave voltage excitation of the magnetic material under test. With the data recorded in this manner Faraday's Law (Ref. 7) is then used to calculate the d\(\phi/\)dt value directly using the excitation voltage level, and the number of turns on the core. In an actual DC-DC converter application the excitation voltage level is the input voltage being applied to a particular magnetic device. Using the SCL characteristics as developed by this method, the power electronics designer may for the first time calculate the SCL for the magnetic device directly knowing only the input voltage, duty cycle, and the frequency of operation.

If the effective cross sectional area (A_e) is known, then the flux density in the magnetic material may be calculated. To accurately determine the value of A_e the physical cross sectional area of the core (A_c) and the stacking factor (SF) must be known. Four methods of

determining the SF value for cut cores are presented in this section.

The results of these methods indicate that the SF value for this

core configuration can be measured to within 5 percent of its actual

value.

The excitation circuit allows three different source configurations to be simulated. By using voltage feedback the excitation source represents a voltage source. If current feedback is used, the excitation source represents a current source. When flux feedback is used, the observed magnetic characteristics become independent of source impedance. Thus, true sinusoidal or square wave voltage excitation can be easily generated. This excitation circuit eliminates many of the problems encountered with conventional square wave excitation sources.

The instrumentation system was designed to efficiently collect the core loss data required for the development of the desired SCL characteristics. The system was also used to run the desired experimental waveshapes against which the analytical core as calculations were then compared. The instrumentation system utilizes a Phillips PM 3252 multiplying oscilloscope for observing instantaneous p(t) waveshapes. A Tektronix 7704 mainframe is used to simultaneously observe e(t), i(t), p(t), $\phi(t)$, and the instantaneous B-H loop.

2.2 Review of Magnetics Equations

All available presentations of specific core loss characteristics display flux density as the independent variable. SCL as the dependent variable, and frequency as a parameter. To use this data requires that the $A_{\bf e}$ of the core be accurately known. This parameter is not always readily available and is inaccurate in most of the published data. However, for the majority of power electronic applications the input

voltage, frequency of operation, number of turns, and duty cycle conditions of a magnetic device are known.

In a magnetic device a time-varying magnetic field produces a voltage which may establish a current in a suitable closed circuit.

Faraday's law (Ref. 7), expressed in the CGS electromagnetic system of units, shows the relationship between this changing magnetic field and the induced voltage:

$$e(\dot{}) = -(d\phi/dt) \times 10^{-8} \text{ volts}$$
 (2.1)

The minus sign indicates that the voltage is in such a direction as to produce a current whose flux, if added to the original flux, would reduce the magnitude of the voltage. This statement that the induced voltage acts to produce an opposing flux is known as "Lenz's law."

If the magnetic field is coupled by an N turn conductor, then

$$e(t) = -N(d\phi/dt) \times 10^{-8} \text{ volts}$$
 (2.2)

By use of equation (2.2), transformer equations for both sinusoidal and quare wave voltage excitation may be derived. If $e(\cdot)$ represents a square wave of voltage, then $\phi(\cdot)$ would be a ramp function with its slope proportional to the volts/turn ratio. Over one half cycle $(0 \le t \le T/2)$, $e(\cdot)$ has a constant value E. For the same interval $\phi(\cdot)$ go s from $-\phi_M$ to $+\phi_M$. The standard form of the transformer equation for square wave excitation is obtained by integrating equation (2.2) over this half cycle interval:

$$\int_{\phi=-\phi_{M}}^{\phi=+\phi_{M}} d\phi = -\frac{E\times10^{8}}{N} \int_{t=0}^{t=T/2} dt$$

$$\phi \bigg]_{-\phi_{M}}^{+\phi_{M}} = -\frac{E \times 10^{8}}{N} t \bigg]_{0}^{T/2}$$

$$2\phi_{M} = -\frac{E \times 10^{8} \text{ T}}{2N}$$

Rearranging terms

$$E = -4\phi_{M}T^{-1} \text{ N} \times 10^{-8} \text{ volts}$$
 (2.3)

Since

$$\phi_{M} = B_{M}A_{e} = B_{M}A_{c}SF \text{ maxwells}$$
 (2.4)

and

$$f = T^{-1} Hz$$
 (2.5)

Substituting (2.4) and (2.5) into (2.3):

$$E = -4B_{M}^{A} \text{ cSFf N} \times 10^{-8} \text{ volts}$$
 (2.6)

If e(t) represents a sinusoidal voltage, then $\phi(t)$ would be a cosine function. Over one half cycle $(0 \le t \le \pi/\omega)$, e(t) is a sine wave with an amplitude of $\sqrt{2}$ E. For the same interval $\phi(t)$ goes from $-\phi_M$ to $+\phi_M$. The standard form of the transformer equation for sinusoidal excitation is obtained by integrating equation (2.2) over this interval:

$$\int_{\phi=-\phi_{M}}^{\phi=+\phi_{M}} d\phi = -\frac{\sqrt{2} E \times 10^{8}}{N} \int_{0}^{\pi/\omega} \sin \omega t \ dt$$

$$\phi = + \frac{\sqrt{2} \times 10^8}{\omega N} \cos \omega t = 0$$

$$2\phi_{M} = -\frac{2\sqrt{2} \times 10^{8}}{\omega N}$$

Rearranging terms

$$E = -(\sqrt{2})^{-1} \phi_{M} \omega \text{ N} \times 10^{-8} \text{ volts}$$
 (2.7)

Since

$$\omega = 2\pi f \text{ rad-sec}^{-1} \tag{2.8}$$

Substituting (2.4) and (2.8) into (2.7)

$$E = -(\sqrt{2})^{-1} B_{MA_{C}}^{A} SF 2\pi i N \times 10^{-8} \text{ volts}$$

or

$$E = -4.44 \, B_{M,C}^{A} \, SF \, f \, n \times 10^{-8} \, volts$$
 (2.9)

By using Equation (2.9) the flux density for a given excitation condition may be calculated. Once this value is known, published SCL data for sinusoidal excitation may be used to directly determine the SCL for a particular magnetic material. To the author's knowledge, published data for square wave excitation conditions is essentially nonexistent. However, using the methods described in this dissertation square wave characteristics may be easily obtained. The squarewave SCL characteristics have a distinct advantage over the sine SCL characteristics because the effects of an additional parameter $d\phi/dt$ may also be examined using the same set of characteristics. In addition, by rearranging the representative axes with frequency as the independent variable, and SCL as the dependent variable, with $B_{\rm M}$ and $d\phi/dt$ as parameters, the characteristics become easier to use. In Chapter 4 the square wave SCL characteristics will be used to gen-

erate the sinusoidal SCL characteristics directly.

The SCL characteristics shown in Appendix A represent the total core loss as a function of frequency. If the energy loss is examined on a per cycle basis (Ref. 8), the effects of $d\phi/dt$, B_M , and frequency may be observed. The energy loss associated with the hysteresis loop is readily determined using Warburg's law (Ref. 9):

$$\tilde{W}_{h} = \frac{1}{4\pi} \iint H dB$$
 watt-sec (2.10)

In the instrumentation system used, the input current to the core under test was observed as the abscissa, and the integral of the input voltage was observed as the ordinate. Using Equation (2.2), and the chain rule, Equation (2.10) may be rewritten in terms of the input voltage $e_i(t)$ and the input current $i_i(t)$:

$$W_{h} = \frac{1}{4\pi} \int_{t=0}^{t=T} H \frac{dB}{dt} \cdot dt$$

$$= -\frac{1}{4\pi} \int_{t=0}^{t=T} \left[NL^{-1} i_{i}(t) \right] \left[A_{e}^{-1} N^{-1} e_{i}(t) \right] dt$$

$$W_{h} = -(4\pi LA_{e})^{-1} \int_{t=0}^{t=T} p_{i}(t) dt \quad \text{watt-sec} \quad (2.11)$$

Equation (2.11) indicates that the area enclosed by the observed B-H loop is proportional to the energy (hysteresis) loss per cycle.

Separation of the eddy current, hysteresis, and anomalous loss components (Refs. 10 to 13) of the total core loss has been explored using only sinusoidal excitation methods. However, by using pulse width modulated square wave excitation techniques a direct separation of loss variables may be accomplished.

2.3 Calculation of Stacking Factor

Four different methods were used to determine the stacking factor of the cut cores used for this dissertation. The oil displacement method and the cross sectional scan method used the physical properties of the core to determine the stacking factor value. The sinusoidal and square wave excitation methods used the electrical properties of the core, and an assumed value for the maximum flux density. Supermalloy cut cores with 1/2, 1, 2, and 4 mil tapes were used and the stacking factors obtained using these four methods were compared with published values. Table I shows the results of this comparison. The four methods used are discussed in detail below:

<u>Displacement method</u>. - In this method Archimedes' principle (Ref. 14) is used to determine the true d. placement volume of the core. First, the core is weighed in air:

$$W_{core} = \rho_{core} V_{Fe}$$

or

TABLE I. - COMPARISON OF FOUR METHODS FOR DETERMINING
STACKING FACTOR WITH PUBLISHED VALUES

Tape thickness, mils	Displacement method	Sine ex. method	Square ex. method	Cross sec. method	Published value
υ.5	0.865	0.885	0.843	0.914	0.80
1.0	.888	.928	.856	.895	.90
2.0	.932	.965	.934	.941	.92
4.0	.927	1.00	.934	.967	.94

Next the core is weighed in an oil of known density. The displacement weight ΔW is then

$$\Delta W = W_{core}$$
 - W_{core} = $\rho_{oil}V_{core}$

or

$$V_{core} = \Delta W(\rho_{cil})^{-1}$$

The stacking factor is the ratio of the air volume of the core to the oil volume of the core:

Stacking factor =
$$\frac{V_{Fe}}{V_{core}}$$

Example:

Core: 1/2 mil Supermalloy cut core $\rho_{\text{core}} = 8.7 \text{ gm-cm}^{-3}$

 $W_{core_{air}}$: 345 grams $\rho_{oil} = 0.928 \text{ gm-cm}^{-3}$

 $W_{core_{oil}}$: 302.5 grams $\Delta W = 42.5$ grams

Stacking factor = (345/8.7)/(42.5/0.928) = 0.865

Excitation methods. - These are indirect methods in which the saturation flux density is used to determine A_e . The core is excited until the current waveform indicates that the core is in the saturation region. Then the excitation voltage level, and frequency are measured. Using the transformer equations derived in Section 2.2, A_e may be calculated:

For the sine wave:

$$A_e = E(4.44 B_M f N \times 10^{-8})^{-1}$$

For the square wave:

$$A_e = E(4B_M f N \times 10^{-8})^{-1}$$

The stacking factor is expressed as

Stacking factor =
$$A_e A_c^{-1}$$

These values are measured at several frequencies and an average stacking factor value is calculated.

Example:

Core:	1/2 mil Supermalloy	A _c :	2.277 cm ²
N: 10	turns	B _M :	7.2 KGauss

For sine wave data:

f, hertz	e, volts	A _e , cm ²
600	3.86	2.01
700	4.48	2.00
800	5.13	2.00
900	5.83	2.02
1000	6.5	2.03

$$\overline{A}_e = 2.01 \text{ cm}^2$$

Stacking factor = 2.01/2.277 = 0.883

For square wave data:

f, hertz	e, volts	A _e , cm ²
600	3.36	1.94
700	3.86	1.91
800	4.41	1.91
900	5.07	1.95
1000	5.44	1.89

$$\overline{A}_e = 1.92 \text{ cm}^2$$

Stacking factor = 1.92/2.277 = 0.843

Cross sectional scan. - In this method a microphotograph of the core cross section is used to determine the stacking factor value. An

optical micrometer is used to measure the adhesive thickness and the actual tape thickness. Since the tape width is constant, the stacking factor is the ratio of the tape thickness to the tape plus adhesive thickness width or

Stacking factor =
$$T_{tape}/(T_{tape} + T_{adhesive})$$

If these values are measured at several points across the core surface, and average stacking factor value may be calculated. In Figure 1 the microphotographs for the four tape thicknesses are shown. Example:

Core: 1/2 mil Supermalloy

Tape thickness, in.	Adhesive thickness, in.
0.000609	0.00003125
.000578	.0000375
.000537	.0000375
•	•
•	•
	• • • • • • • • • • • • • • • • • • • •
0.000546	0.0000468
$\overline{T}_{tape} = 0.000546 \text{ (in.)}$	$\bar{T}_{adhesive} = 0.00005125 (in.)$
Stacking factor = 0.000	546/(0.000546 + 0.00005125)

2.4 General Experimental Set-Up

= 0.914

Previous excitation methods have used feedback control schemes

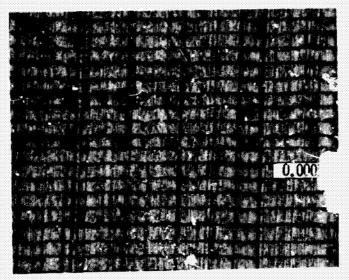
(Ref. 15) to reduce the harmonic content of the sinusoidal excitation

by controlling the apparent source impedance. For this work, feedback

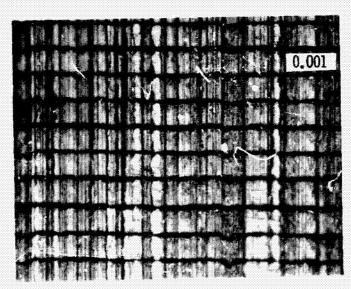
control is used to achieve the desired voltage excitation. If a square

wave voltage is required, a triangular reference voltage is used. If

an impulse voltage is required, a square wave reference voltage is used.

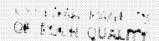


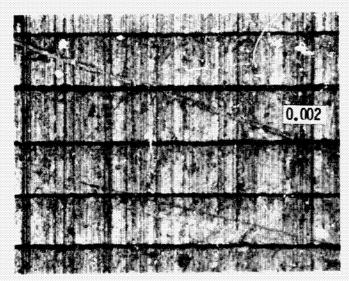
(a) 1/2 mil Supermalloy tape.



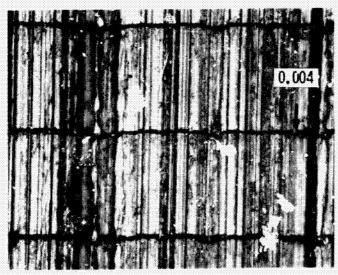
(b) 1 mil Supermallov tape.

Figure 1. - Cross section Microphotogra- of 1/2, 1, and 4 mil Supermalloy cut cores.





(c) 2 mil Supermalloy tape.



(d) 4 mil Supermalloy tape.

Figure 1. - Concluded.

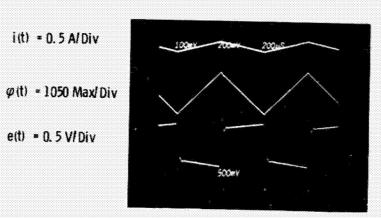


And, if sinusoidal voltage excitation is required, a cosinusoidal reference is used. Figure 2 shows the voltage, current, and flux relationsnips for each of these flux reference conditions.

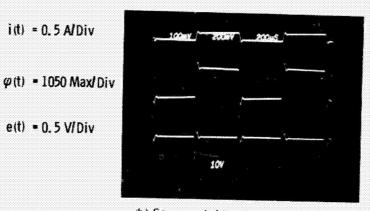
The complete instrumentation system is shown in Figure 3. A high power operational amplifier was designed which allowed the option of selecting eitner voltage, current, or the integral of voltage (flux) as the feedback parameter. A Wavetek model 146 signal generator was used to generate the desired symmetrical flux waveshape. An ICL 8038 multifunction integrated circuit was used to generate the asymmetrical square and sinusoidal reference waveforms. The excitation frequency was measured with a Hewlett-Packard 5212A electronic counter. A Clarke-Hess model 225 wideband wattmeter (0 to 300 kHz) was used to measure the average input power into the core under test. This instrument allows the power delivered by a 100 kHz square wave to be normally measured to within 2 percent. The internal current shunt in the wattmeter displays the input current signal used for several measurements. A Phillips PM 3252 multiplying oscilloscope was also used to obtain the instantaneous $p_{i}(t)$ in the core. The input voltage to the core was measured with a Hewlett-Packard 3403C true RMS voltmeter. The Tektronix 7704 oscilloscope was used to observe $i_{i}(t)$, $\mathbf{e_i}(t)$, $\phi(t)$, $\mathbf{p_i}(t)$ as well as the instantaneous B-H loop of the core under test.

2.5 Source Impedance Effects

Source impedance effects discussed in the literature (Refs. 16 and 17) show that the width of the B-H loop to be dependent upon the driving source impedance. By changing the feedback parameter on the operational amplifier, a voltage. $\langle R_{\rm T} \rangle$ or a current source



(a) Ramp døldt - square voltage.



(b) Square $d\phi / dt \rightarrow impulse voltage$.

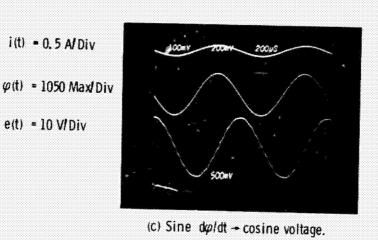


Figure 2. - Comparisons of i(t), $\varphi(t)$, and e(t) for ramp, square, and cosinusoidal reference voltages.



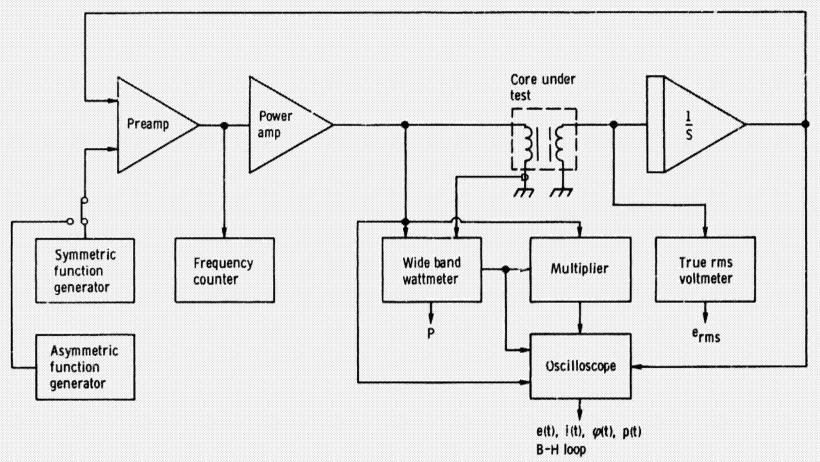


Figure 3. - Detailed instrumentation system.

 $(R_S >> R_L)$ may be simulated. The flux, current, and B-H loop characteristics for these two source configurations are snown in Figures 4 and 5. For the same B_M conditions, the width of the B-H loop is wider if the core is being driven by a current source. However, if the flux is observed for these two source conditions it can be seen that for sinusoidal current conditions, the flux approaches a square wave. This is a condition having a higher $d\phi/dt$ rate in the core than that realized under sinusoidal voltage conditions. Therefore, the observed B-H width is really a function of $d\phi/dt$ rather than being dependent upon source impedance.

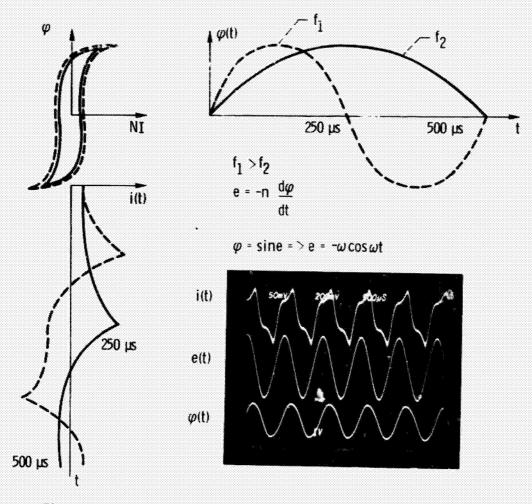


Figure 4. – Relationships for i(t), φ (t), and φ -I under voltage excitation.

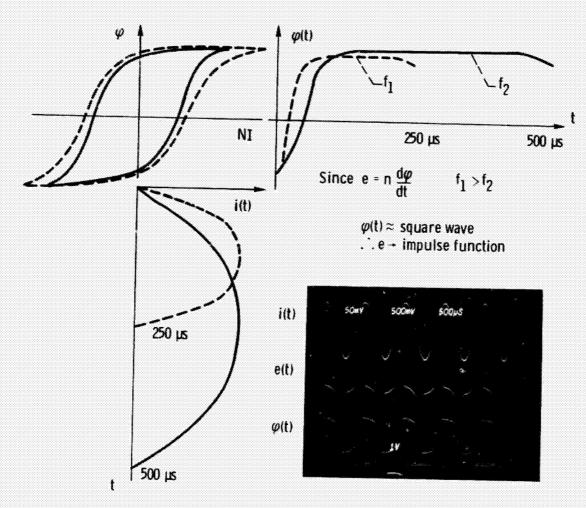


Figure 5. - Relationships for i(t), $\varphi(t)$, and φ -I under current excitation.

CHAPTER III

ANALYTICAL CORE LOSS CALCULATIONS

3.1 Introduction

Analytical prediction of core loss in magnetic materials is important from both an economic (Ref. 18) and an engineering design (Ref. 19) standpoint. Accurate prediction allows design tradeoffs to be made between the areas of weight, efficiency, and performance. In addition new circuit developments employing magnetic devices may be more fully analyzed without hardware fabrication. Early work (Ref. 20) at frequencies of up to 200 kHz was aimed at measuring the magnetic properties of iron. Separation of core loss into its specific components of eddy current, static hysteresis, and anomalous loss has been examined in great depth (Refs. 21 to 23) for low frequency (<1 kHz) sinusoidal excitations.

With the growth of power electronics to higher power levels and the employment of new converter techniques, the need to understand the properties of magnetic materials under nonsinusoidal conditions is arising. Recent comparisons between sinusoidal and square wave core loss (Ref. 5) indicate a lack of theory to predict the loss difference observed between these two excitations. However, if the dependence of the B-H characteristics upon the flux conditions described in Section 2.5 is expanded, the loss characteristics for these two excitations may be qualitatively observed, and predicted using easily obtained data.

In this chapter a comparison between sinusoidal and square wave

core loss is made. The results are shown to be consistent with the values published in the literature. Next, pulse width modulated square wave excitation is used as an analytical tool to achieve a true severation between the $d\phi/dt$ and B_M effects. Finally, three methods are developed to analytically predict the SCL for square wave excitation. These results are compared with experimentally measured results obtained under the same excitation conditions.

3.2 Specific Core Loss Measurements

To generate the SCL characteristics under square wave excitation for the magnetic materials studied, data was taken by holding $d\phi/dt$ constant and varying the excitation frequency. The power input to the core was measured at regular frequency intervals. The actual values for SCL in walts/pound and $B_{\rm M}$ in gauss were calculated from the raw data using the following formulas:

Measured e_{rms} in volts

$$\frac{1}{kA_c SFf N \times 10^{-8}} = B_M (gauss)$$
 (3.2)

where

k = 4 for square wave excitation

= 4.44 for sinusoidal excitation

Figure 6 shows the specific core loss for 1/2 mil cut Supermalloy under both sinusoidal (Fig. 6(a)) and square wave (Fig. 6(b)). In addition, comparisons are made with published sinusoidal data for these two excitation conditions. The major difference between the published and measured values for sinusoidal excitation may be accounted for by

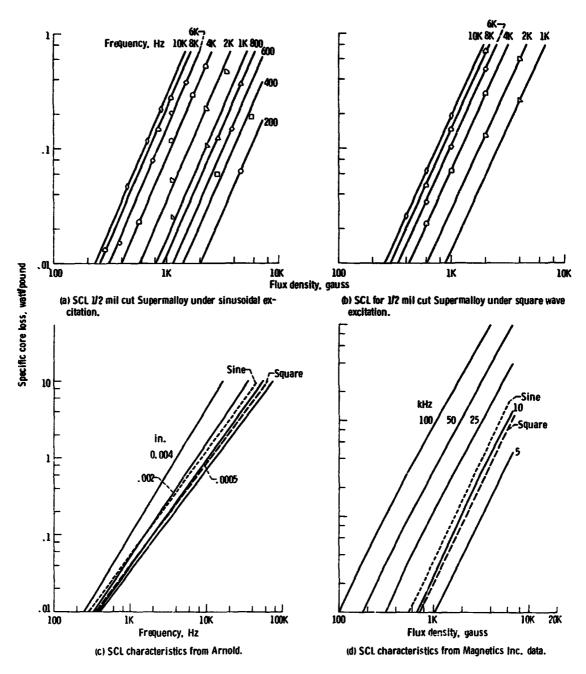


Figure 6. - Comparison of SCL characteristics for sinusoidal and square wave excitation with published data.

several factors:

Probable accuracy of the published data

Accuracy of the core loss measurement

Limited raw core loss data available

Driving source impedance

Measurement of actual excitation voltage level

Value assumed for stacking factor

Batch variations for the magnetic material

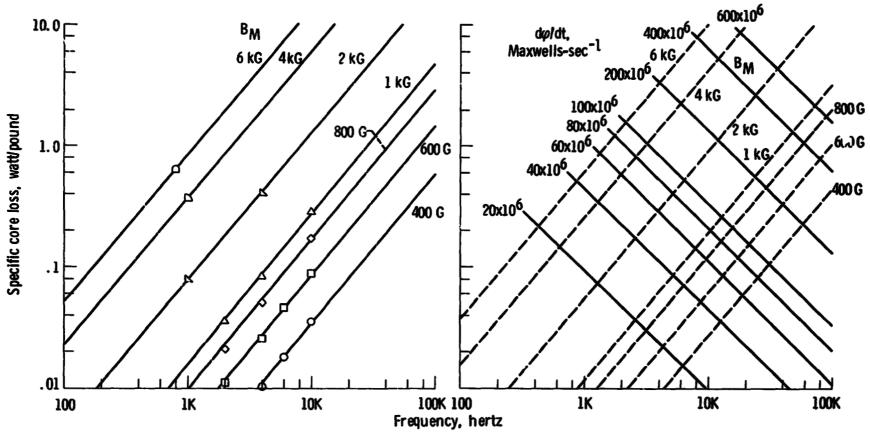
Actual lamination thickness

It is interesting to note that for equal values of $B_{\underline{M}}$ the SCL for sinusoidal excitation is greater than for square wave excitation. Example:

B_M = 2k gauss frequency = 1 kHz SCL (sine) = 0.078 W/1b SCL (square) = 0.058 W/1b

From these results it can be seen that the SCL for square wave excitation is approximately 26 percent less than the SCL for sinus-oidal excitation. This result is consistent with previous observations (Ref. 5) which compared loss conditions for these two excitations.

To aid in the development of methods for specific loss calculations, data is replotted with the frequency and flux density axes interchanged. Figure 7(a) shows SCL versus frequency with $B_{\rm M}$ as a parameter for cut 1/2 mil Supermalloy under sinusoidal excitation. Figure 7(b) shows these characteristics for square wave excitation. For square wave excitation conditions, $d\phi/dt$ was held constant as well as frequency. This allowed the additional parameter of $d\phi/dt$ to be



- (a) Specific core loss characteristics for cut 1/2 mil Supermalloy for sinusoidal excitation.
- (b) Specific core loss characteristics for cut 1/2 mil Supermalloy for square wave excitation with B_M and døldt as parameters.

Figure 7. - SCL characteristics for sinusoidal and square wave excitation with axes interchanged.

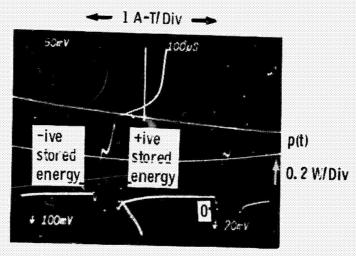
plotted on the same SCL characteristic (Fig. 7(b)).

3.3 Separation of Loss Variables at Constant Frequency

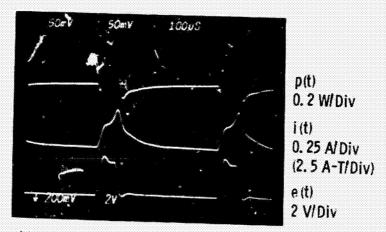
If excitation is restricted to that of a constant frequency pulse width modulated square wave, the $d\phi/dt$ effects may be determined for various constant values of B_M . Figure 8 shows the instantaneous waveforms for p(t), i(t), and e(t). Positive i(t) values correspond to $H(Ni(t)L^{-1})$ values to the left of the B-axis while negative i(t) values correspond to H values on the right of the B-axis. Also, p(t) values are both positive and negative which corresponds to energy supplied to or energy returned from the core.

Figure 9 shows the instantaneous B-H loop of 1/2 mil Supermalloy measured for various $d\phi/dt$ (change in duty cycle) values while holding B_M constant at 6.76 K gauss, and frequency constant at 2 kHz. From this figure, it may be seen that as $d\phi/dt$ increases the B-H loop area for that portion of the cycle also increases. Since the B-H area is proportional to the energy loss, and the additional B-H area is proportional to $d\phi/dt$, it is assumed that the additional half cycle energy loss observed is proportional to $d\phi/dt$.

This procedure was then repeated for several values of B_M with frequency held constant with the results shown in Figure 10. From this figure it is seen that for constant $d\phi/dt$, the total core loss increases with B_M . If the lines of constant B_M could be linearly extrapolated, the ordinate $(d\phi/dt = 0)$ intercept would indicate the hysteresis increase with B_M . The eddy current loss is dependent upon both $d\phi/dt$ and B_M . However, it is interesting to note that for a constant value of B_M , the half cycle loss is linearly dependent on



(a) B-H loop and instantaneous p(t).



(b) Instantaneous p(t), i(t), and e(t).

Figure 8. - Comparison of the four loss regions for the instantaneous B-H loop, p(t), i(t), and e(t).



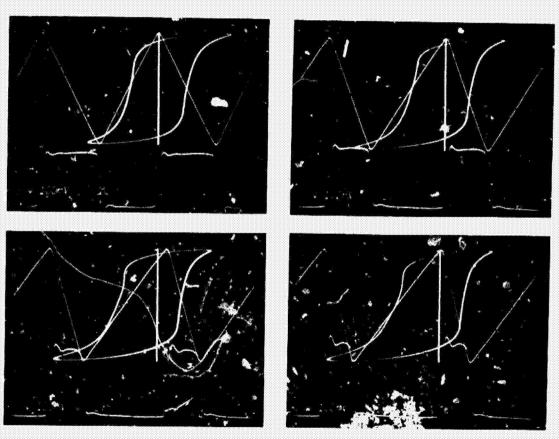


Figure 9. B-H loop as a function of $\,\mathrm{d}\varphi\mathrm{Idt}\,$ for constant B_m and f.

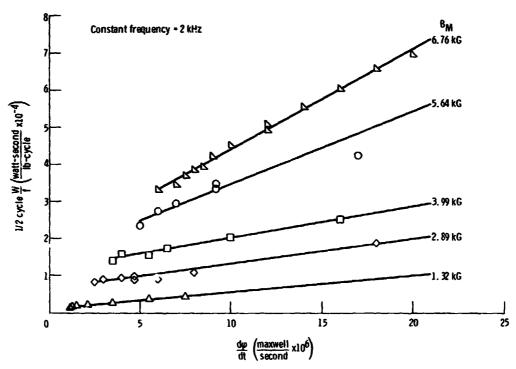
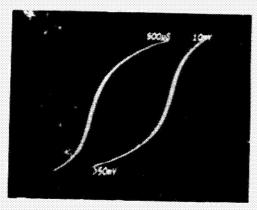


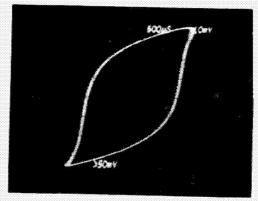
Figure 10. – Half cycle specific core loss per cycle vs. d ϕ ldt with $\,$ B $_{\mbox{\scriptsize M}}\,$ as a parameter.

 $d\phi/dt$. This procedure shows that a true separation of the variables $d\phi/dt$, and B_M may be achieved using pulse width modulated square wave excitation techniques. These characteristics will be used to qualitatively explain the differences observed in the B-H loop for sinusoidal and square wave excitation.

A comparison of the B-H loop characteristics under both constant (square wave (fig. 11(a)) and variable (sinusoidal (fig. 11(b)) $d\phi/dt$ conditions measured while holding both frequency and B_M (relates to hysteresis loss) constant is shown in Figure 11. Figure 11(c) illustrates some interesting conditions which, to the author's knowledge, have never been shown in the literature. It has been shown (Fig. 10) that the total half cycle core loss is composed of a hysteresis loss which is proportional to B_M , and an induced eddy current loss proportional to both $d\phi/dt$ and B_M .

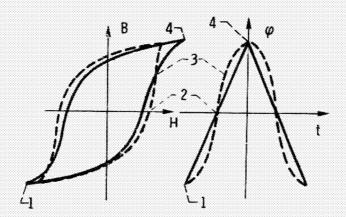
Since B_M is held constant, hysteresis loss is constant, and eddy current loss is dependent on $d\phi/dt$ only. The following qualitative observations may be made. At point 1, $d\phi/dt$ for the sine wave is equal to zero while $d\phi/dt$ for the square wave is equal to some positive constant value greater than zero. Therefore, the core loss for the square wave is greater (wider B-H width) than that for the sine wave at point 1. As the flux density increases to point 2, $d\phi/dt$ for the sine wave reaches its maximum value. This value of $d\phi/dt$ is reater than the $d\phi/dt$ for the square wave, and the corresponding core loss (B-H width) is greater. As the flux density increases to point 3, $d\phi/dt$'s for the sine and square wave are approximately equal. This corresponds to an equal B-H width. As the flux approaches point 4 the same conditions as point 1 are repeated.





(a) B-H loop with $d\varphi/dt = k$

(b) B-H loop with $d\varphi/dt = k(t)$.



(c) Superimposed B-H loops and flux excitations for constant and variable $\mathrm{d}\varphi/\mathrm{d}t$ conditions.

Figure 11. - Comparison of B-H loop characteristics for constant and variable ${\rm d}\varphi/{\rm d}t$ conditions with frequency and ${\rm B_m}$ as a constant.

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3.4 Analytical Calculation Methods

Using the SCL characteristics for square wave excitation three methods were developed which analytically predict the SCL of a magnetic material for pulse width modulated square wave excitation. These methods are discussed in detail below.

Method 1 - Specific core loss based on $B_{\rm M}$ of the core. - The SCL characteristics presented in Figure 7(b) relate the total specific loss (hysteresis + eddy current = total loss) to frequency with $B_{\rm M}$ as a parameter. This data will be used to determine the core loss characteristics with pulse width modulated square wave core excitation. This method is the most indirect approach to calculating the specific core loss. It requires an accurate measure of the effective cross sectional area of the core loss characteristics with pulse with modulated square wave core excitation. This method is the most indirect approach to calculating the specific core loss. It requires an accurate measure of the effective cross sectional area of the core. Generally, this is a parameter which is not accurately k. the power electronics designer.

Figure 1 shows the oscilloscope traces of a 1000 Hz square wave operating over a duty cycle range of 10 to 50 percent. For each duty cycle the voltage for the positive cycle is adjusted to a maximum value of 4 volts. The DC offset in the core is held to a minimum. Therefore, the negative and positive half cycle volt-seconds are equal.

In order to use the Figure 7(b) characteristics the $\,^{}{\rm B}_{\rm M}\,^{}$ value for each duty cycle condition must be determined. This value may be obtained as follows:

3.4 Analytical Calculation Methods

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In order to use the Figure 7(b) characteristics the $\, B_{M}^{} \,$ value for each duty cycle condition must be determined. This value may be obtained as follows:

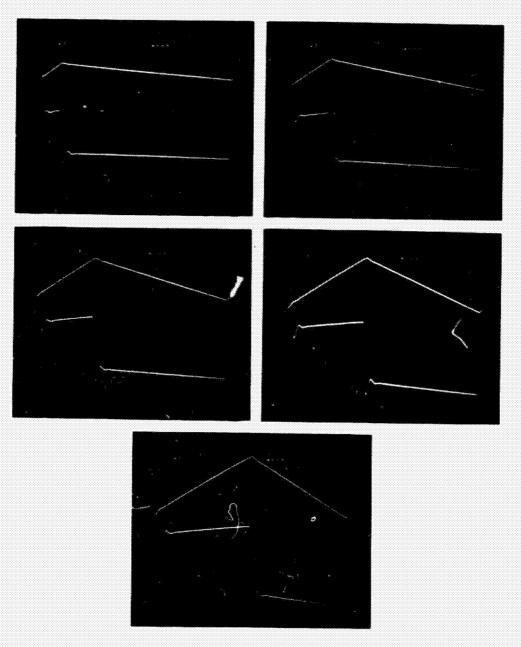


Figure 12. - Waveform excitation for a pulse width modulated square wave.

ORIGINAL PAGE IS OF POOR QUALTE From Faraday's law

$$e_c = -N d\phi/dt \times 10^{-8} \text{ (volts)}$$
 (3.3)

Therefore

$$d\phi/dt = -e_c N^{-1} \times 10^8$$
 (maxwells/second) (3.4)

For each cycle portion

$$\Delta \phi = (d\phi/dt)\Delta t \tag{3.5}$$

$$\Delta \phi = -e_c N^{-1} \Delta t \times 10^8 \tag{3.6}$$

and,

$$B_{M} = \Delta \dot{\phi} (2A_{e})^{-1} = -0.5 e_{c} \Delta t (NA_{e})^{-1} \times 10^{8}$$
 (gauss) (3.7)

where:

$$A_e = A_c \times SF (cm^2)$$

N = number of turns

 e_c , $\Delta t \rightarrow related to duty cycle$

Using 1/2 mil Supermalloy as an example, for a frequency of 1000 Hz, the B_{M} values corresponding to each duty cycle width can be determined as follows:

$$B_{M}$$
 (positive cycle) = -0.5 (3.9) (100×10⁻⁶) (10×1.97)⁻¹×10⁸
= -990 gauss

$$B_{M}$$
 (negative cycle) = -0.5 (-0.44) (900×10⁻⁶) (10×1.97)⁻¹×10⁸
= 1000 gauss

$$\mathbf{B}_{\mathbf{M}}$$
 (positive cycle) $\mathbf{\tilde{a}}$ $\mathbf{B}_{\mathbf{M}}$ (negative cycle)

Considering the positive half cycle only this calculation can be repeated for the remaining four cases with the results shown below:

Duty cycle, percent	Δt−μs	e _c , volts	B _M ,
10	100	3.9	990
20	200	3.9	1980
30	300	3.9	2470
40	400	3.8	3858
50	500	3.4	4315

This data can now be used with Figure 7(b) to determine the specific core loss associated with each duty cycle condition. The results of these calculations are shown on Figure 13.

During the experiment the specific core loss was measured for each duty-cycle condition. A comparison between the measured and calculated results is shown in Figure 14. At the 50 percent duty-cycle condition the error between the measured and calculated specific core loss was:

$$SCL_{calc} = 0.31 \text{ W/lb}$$

$$SCL_{meas} = 0.339 \text{ W/lb}$$

percent error =
$$\frac{SCL_{meas} - SCL_{calc}}{SCL_{meas}} \times 100$$
 (3.8)

percent error =
$$\frac{0.339 - 0.31}{0.339} \times 100 = 8.6$$
 percent

From these results it may be seen that the specific core loss can be easily calculated for pulse width modulated square waves if the base frequency, excitation voltage level, duty-cycle, number of turns, and the effective cross sectional area of the core is known. Also, for a constant input voltage as the duty cycle increases, $\mathbf{B}_{\mathbf{M}}$ increases, and the specific core loss of the magnetic device increases.

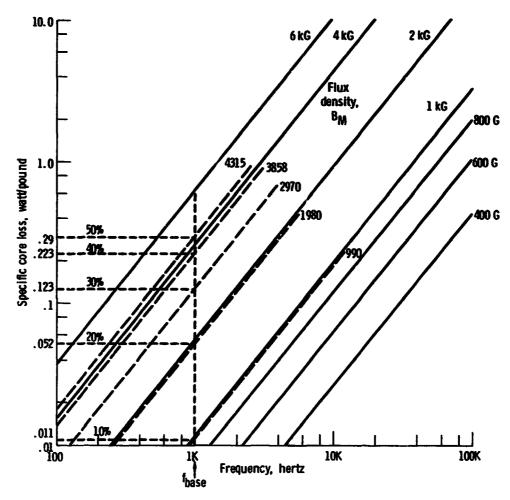


Figure 13. - Specific core loss characteristics for cut 1/2 mil Supermalloy under duty cycle conditions.

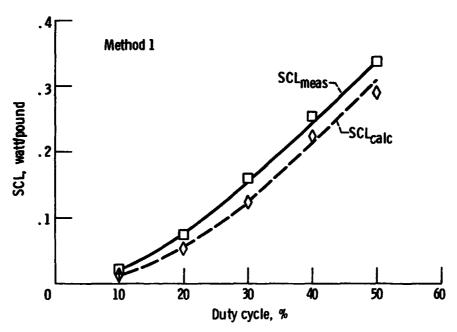


Figure 14. - Comparison of calculated vs measured values of SCL for 1/2 mil Supermalloy as a function of duty cycle for a 1 kHz - 3.9 volt pulse voltage.

Method 2 - Specific core loss based on $d\phi/dt$ in the core. - The specific core loss characteristics presented in Figure 7(b) relate the total specific core loss to frequency with B_M and $d\phi/dt$ as parameters. If the B_M parameter is ignored on this figure, the specific core loss can be calculated by using the $d\phi/dt$ parameter and an equivalent excitation frequency.

In the discussion of Section 3.3 the core loss was shown to be a function of B_M and $d\phi/dt$. For a pulse width modulated square wave with a constant positive $d\phi/dt$ and constant negative $d\phi/dt$, B_M is determined directly by the length of time that the pulse is on. The width of this "on pulse" has an "equivalent" frequency of

$$f_{eq} = T_{eq}^{-1} = (2t_{on})^{-1}$$
 (3.9)

Also, it was shown that the total specific core loss per cycle could be broken down into a half cycle loss dependent on $d\phi/dt$. Therefore, the total specific core loss may be determined using the following relationship.

Total SCL = (SCL/cycle) (excitation frequency)

$$= \left(\frac{1}{2} \frac{\text{SCL}^{+}}{f_{\text{eq}}^{+}} + \frac{1}{2} \frac{\text{SCL}^{-}}{f_{\text{eq}}^{-}}\right) f_{\text{ex}}$$
 (3.10)

where

SCL⁺ positive cycle specific core loss

SCL negative cycle specific core loss

 f_{eq}^{+} positive cycle equivalent frequency

f negative cycle equivalent frequency

f excitation frequency

Example:

10 percent duty-cycle case:

$$e_{c}^{+} = 3.9 \text{ V}$$

$$e_{c}^{-} = 0.44 \text{ V}$$

$$f_{ex} = 1000 \text{ Hz}$$

$$t_{on}^{+} = 100 \times 10^{-6} \text{ sec}$$

$$t_{an}^{-} = 900 \times 10^{-6} \text{ sec}$$

Then

$$f_{eq}^{+} = (2t_{on}^{+})^{-1} = (2 \times 100 \times 10^{-6})^{-1} = 5 \text{ kHz}$$

$$f_{eq}^- = (2t_{on}^-)^{-1} = (2 \times 900 \times 10^{-6})^{-1} = 555 \text{ Hz}$$

using Faraday's law

$$d\phi/dt = e_c N^{-1} \times 10^8$$

$$(d\phi/dt)^+ = 3.9(10)^{-1} \times 10^8 = 39 \times 10^6$$
 maxwells/sec

$$(d\phi/dt)^- = 0.44(10)^{-1} \times 10^8 = 4.4 \times 10^6$$
 maxwells/sec

Replotting Figure 7(b) as Figure 15 with $(d\phi/dt)^+$, $(d\phi/dt)^-$, f^+ , and f^- shown, the corresponding SCL's may be read directly from the figure:

Using the equation for the total SCL the following results are obtained:

Total SCL =
$$\frac{1}{2} \left(\frac{0.082}{5000} + \frac{0.0063}{555} \right) 1000$$

Total SCL = 0.014 watt/pound

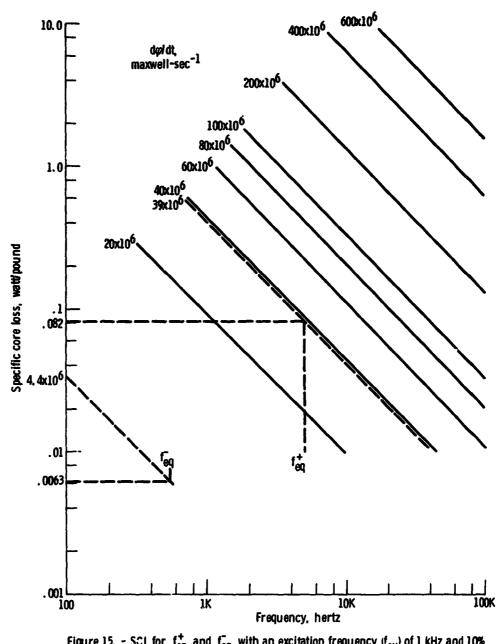


Figure 15. - SCL for f_{eq}^+ and f_{eq}^- with an excitation frequency (f_{ex}) of 1 kHz and 10% duty cycle.

summarizing the results for the other conditions we have:

Duty cycle, percent	f ⁺ eq, Hz	f ⁻ eq, Hz	$\frac{d\phi^{+}}{dt}$ ×10 ⁶ max/sec	$\frac{d\phi}{dt}$ ×10 ⁶ max/sec
10	5000	555	39	4.4
20	2500	625	39	10
30	1670	714	39	16
40	1250	833	38	26
50	1000	1000	34	34

From Figure 15 the SCL⁺ and SCL⁻ can be obtained. Again using the total SCL equation:

20 percent duty cycle:

Total SCL =
$$\frac{1}{2} \left(\frac{0.162}{2500} + \frac{0.0275}{625} \right) 1000 = 0.054 \text{ W/lb}$$

30 percent duty cycle:

Total SCL =
$$\frac{1}{2} \left(\frac{0.24}{1670} + \frac{0.074}{714} \right) 1000 = 0.124 \text{ W/lb}$$

40 percent duty cycle:

Total SCL =
$$\frac{1}{2} \left(\frac{0.3}{1250} + \frac{0.192}{833} \right) 1000 = 0.235 \text{ W/lb}$$

50 percent duty cycle:

Total SCL =
$$\frac{1}{2} \left(\frac{0.3}{1000} + \frac{0.3}{1000} \right) 1000 = 0.3 \text{ W/lb}$$

In Figure 16 these calculated SCL's are compared with the measured SCL values for these duty-cycle conditions. The percent error is the same as obtained with method 1 (8.6 percent).

Method 3 - Specific core loss based on incremental $d\phi/dt$ in the core. - The technique developed in method 2 allows determination of the specific core loss based on the $d\phi/dt$ value for the entire positive or the entire negative cycle of the square wave. In Sec-

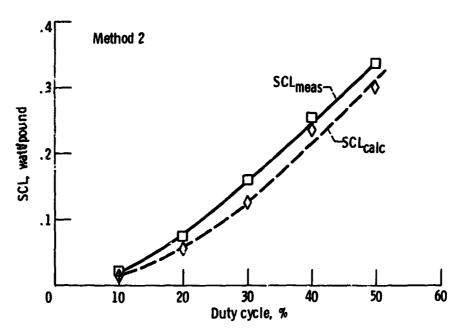


Figure 16. - Comparison of calculated vs measured values of SCL as a function of duty-cycle for a 1 kHz - 3. 9 volt pulse voltage.

tion 3.3 the dependency of the B-H width on $d\phi/dt$ was developed. For a particular frequency and B_M level a family of curves could be developed for various $d\phi/dt$ values as shown in Figure 17. Using actual square waves observed in this experiment (Fig. 12), it can be seen that for any interval the actual exciting voltage is not a constraint value. This would indicate that the corresponding $d\phi/dt$ values over the cycle are not of constant value. If the waveform is broken down into discrete sample intervals the corresponding locus of points on the B-H loop may be plotted as Δt is incremented. The corresponding B-H width would be a direct function of the pulse height for that interval. Figure 17(b) shows how the pulse height, $d\phi/dt$, and Δt would translate into a locus of points on the B-H loop. This approach may be used to explain the B-H loop shapes observed under various excitation conditions.

Relating these various $d\phi/dt$ rates to loss at the positive equivalent frequency, the SCL characteristics of Figure 17(c) may be used to determine the incremental core loss.

The loss per interval is calculated as follows:

$$T^{+} = n \Delta t_{i}^{+} = \frac{1}{2f^{+}}$$
 (3.11)

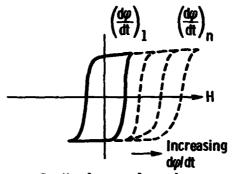
where n = number of increments

$$\Delta t_{i}^{+} = \frac{1}{2nf^{+}} \tag{3.12}$$

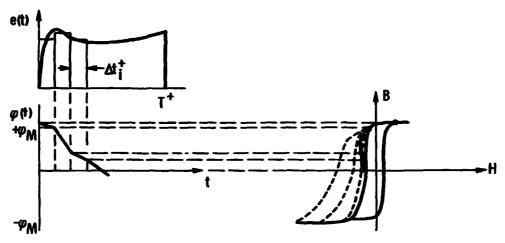
The energy dissipated in Δt_i^+ would be:

$$\Delta E_{i}^{\dagger} = SCL_{i}^{\dagger} \Delta t_{i}^{\dagger} = \frac{SCL_{i}^{\dagger}}{2nf^{\dagger}}$$
 (3.13)

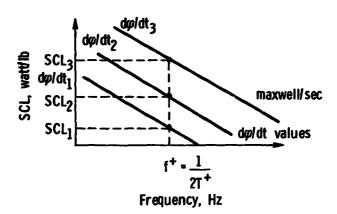
If Δt_i^+ were made small enough, the difference between the incre-



(a) Family of curves for various doldt values.



(b) Relationship of incremental doldt rates to B-H loop.



(c) SCL characteristics for various døldt values.

Figure 17. – Family of B–H loops and SCL characteristics based on $\,\mathrm{d} \varphi \mathrm{I} \,\mathrm{d} t\,$ as a parameter.

mental excitation and the actual excitation would approach zero as a limit.

The total energy over the positive cycle is simply the summation of the incremental ΔE^{+} s:

$$E_{T}^{+} = \underbrace{\sum_{i=1}^{N} \Delta E_{i}^{+}}_{i} = \underbrace{\frac{1}{2nf^{+}} \sum_{i=1}^{n} SCL_{i}^{+}}_{i} \underbrace{W-sec}_{lb-cycle}$$
(3.14)

Similarly the total energy for the negative cycle is

$$E_{T}^{-} = \underbrace{\frac{M}{\sum_{i=1}^{M}} \Delta E_{i}^{-}} = \underbrace{\frac{1}{2mf}}_{i=1}^{m} SCL_{i}^{-} \underbrace{\frac{W-sec}{lb-cycle}}_{1b-cycle}$$
(3.15)

and the average power would be:

$$SCL_{T} = specific power loss = \frac{E_{T}^{+} + E_{T}^{-}}{T^{+} + T^{-}} \qquad \text{W/1b} \qquad (3.16)$$

Applying this technique to the 10 percent PWM square wave case we have:

Positive cycle

Interval $(\Delta t_i = 20 \mu s)$	e _c , volts	$\frac{d\phi}{dt} \div 10^6$ max/sec	SCL ⁺ , W/lb
1	3.4	34	0.063
2	4.0	40	.092
3	3.9	39	.087
4	3.9	39	.087
5	4.0	40	092
			Σ 0.421

Negative cycle

Interval $(\Delta t_i = 100 \mu s)$	e _c , volts	$\frac{d\phi}{dt}$: 10^6 max/sec	SCL _i , W/1b
1	0.2	2	0.00092
2	.4	4	.0044
3	.4	4	.0044
4	.4	4	.0044
5	.44	4.4	.0055
6	.5	5	.0074
7	.5	5	.0074
8	.55	5.5	.0091
9	.55	5.5	.0091

Σ 0.05262

$$E_{T}^{+} = \frac{1}{2.\times 5 \times 5000} (0.421) = 8.42 \times 10^{-6} \frac{\text{W-sec}}{\text{lb-cycle}}$$

$$E_{T}^{-} = \frac{1}{2 \times 9 \times 555} (0.05262) = 5.267 \times 10^{-6} \frac{\text{W-sec}}{\text{lb-cycle}}$$

$$SCL_{T} = \frac{8.42 \times 10^{-6} + 5.267 \times 10^{-6}}{100 \times 10^{-6} + 900 \times 10^{-6}} = 0.014 \text{ W/lb}$$

For the remaining cases we have

20 percent:

$$\sum_{i=1}^{10} SCL_{i}^{+} = 1.764 \qquad \Delta t_{i}^{+} = 20 \ \mu s$$

$$\sum_{i=1}^{8} SCL_{i}^{-} = 0.298 \qquad \Delta t_{i}^{-} = 100 \ \mu s$$

$$SCL_{T} = \frac{\frac{1.764}{2 \times 10 \times 2500} + \frac{0.298}{2 \times 8 \times 625}}{200 \times 10^{-6} + 800 \times 10^{-6}} = 0.065 \text{ W/1b}$$

30 percent:

$$\sum_{i=1}^{15} SCL_{i}^{+} = 4.554 \qquad \Delta t_{i}^{+} = 20 \text{ } \mu s$$

$$\sum_{i=1}^{7} SCL_{i}^{-} = 0.621 \qquad \Delta t_{i}^{-} = 100 \text{ } \mu s$$

$$SCL_{T} = \frac{4.554}{2 \times 15 \times 1670} + \frac{0.621}{2 \times 7 \times 714} = 0.153 \text{ W/lb}$$

40 percent:

$$\sum_{i=1}^{20} SCL_{i}^{+} = 8.204 \qquad \Delta t_{i}^{+} = 20 \text{ } \mu s$$

$$\sum_{i=1}^{6} SCL_{i}^{-} = 1.117 \qquad \Delta t_{i}^{-} = 100 \text{ } \mu s$$

$$SCL_{T} = \frac{7.73}{\frac{2 \times 20 \times 1250}{400 \times 10^{-6} + 600 \times 10^{-6}}} = 0.266 \text{ W/lb}$$

50 percent:

$$\sum_{i=1}^{25} SCL_{i}^{+} = 9.03 \qquad \Delta t_{i}^{+} = 20 \text{ µs}$$

$$\sum_{i=1}^{5} SCL_{i}^{-} = 1.66 \qquad \Delta t_{i}^{-} = 100 \text{ µs}$$

$$SCL_{T} = \frac{9.03}{\frac{2 \times 25 \times 1000}{500 \times 10^{-6} + 500 \times 10^{-6}}} = 0.347 \text{ W/1b}$$

A comparison of this value with the measured results is shown in

Figure 18. The error using this technique is less than that of methods 1 and 2.

The instantaneous p(t) waveforms are shown in figure 19. A

Phillips PM3252 multiplying oscilloscope was used to obtain p(t). In

the top photo a resistive load was used to verify the accuracy of the

multiplier. For the square wave the growth of p(t) is fairly linear

with time, which supports the preliminary assumptions used in developing this method.

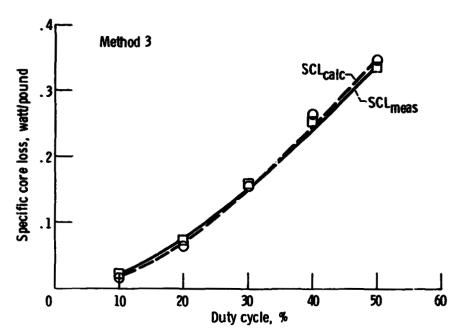


Figure 18. - Comparison of calculated vs measured values of SCL as a function of duty cycle for a 1 kHz - 3.9 volt pulse voltage.

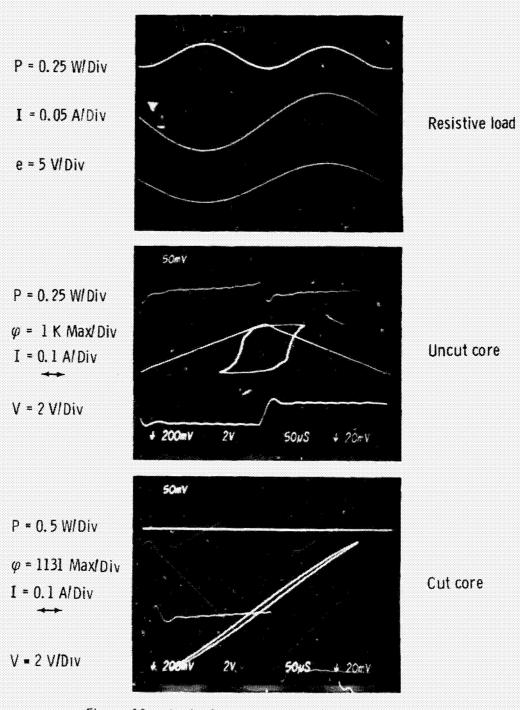


Figure 19. - Instantaneous energy characteristics.

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CHAPTER IV

GENERAL CORE LOSS ANALYSIS

4.1 Introduction

Analysis techniques such as harmonic analysis and digital simulation of eddy current losses have previously been used for prediction of core losses of distorted flux waveforms (Ref. 24). Eddy current correction factors (Ref. 25) have also been employed in estimating the core loss of sinusoidal excitation considering only the fundamental and one harmonic component. In this chapter method 3 of the previous chapter is used to analytically predict the specific core loss for several distorted waveform excitations.

Since this method is based upon the summation of a large number of sample points, the digital computer was used for high speed core loss analysis. In order to minimize the number of subroutines of the computer algorithm, a general characteristic equation was derived using experimental data. This equation relates SCL to the equivalent frequency and $d\phi/dt$. An alternate equation which relates SCL to equivalent frequency and B_{M} was also derived. For analysis of the SCL characteristics a specific magnetic material having a known tape or grain thicknesses, only a constant and two exponent values must be supplied in the initial data format.

Using the square wave SCL characteristics and the computer algorithm, analyses were performed for sinusoidal and nonsinusoidal excitation conditions using Supermalloy, Square Permalloy 80, and ferrite

materials. As a final illustration of the method, the sinusoidal core loss characteristics of a 1/2 mil Supermalloy material operating at 1000 Hz is measured experimentally and calculated analytically.

4.2 Development of Computer Algorithms

The SCL characteristics for various types of magnetic materials is presented in appendix A. Due to the logarithmic nature of the characteristics, the SCL equation for either the $d\phi/dt$ parameter or the B_M parameter would have the following form:

$$SCL = af_{eq}^{X}$$
 (4.1)

The exponent x could be either positive (let x = d) or negative (let x = u) depending upon the parameter being examined. The constant a is also a logarithmic function of either $d\phi/dt$ or B_M . The equation for the constant is as follows:

$$a = g(d\phi/dt)^{c}$$
 (for the $d\phi/dt$ parameter) (4.2)

or

$$a = s(B_M)^Z$$
 (for the B_M parameter) (4.3)

Combining Equation (4.1) with the appropriate constant equation, the following two equation forms may be used to determine the SCL for a specific frequency and B_M or $d\phi/dt$ level:

For d\psi/dt:

$$SCL = g(d\phi/dt)^{c} f_{eq}^{u}$$
 (4.4)

For B_{M} :

$$SCL = s(B_{M})^{z} f_{eq}^{d}$$
 (4.5)

Values for the constants g and s, and for the exponents c, d, z, and μ for various magnetic materials were determined from the

experimental data shown in appendix A. These calculated values are shown in tables II and III. In Equation (4.5) at low flux densities, the d exponent was determined by Steinmetz to be 1.6 for electrical sheet steel (Ref. 26). For higher flux densities the value of the exponent may be as high as 2.5. The values of the exponent as measured in these experiments falls within this upper limit. The value of the exponent was found to be dependent upon both the magnetic material and the tape thickness.

TABLE II. - TABULATION OF THE CONSTANT AND EXPONENT

VALUES USED IN EQUATION (4.4)

Table	II.1.	_	Supermalloy-Cut	Core
-------	-------	---	-----------------	------

Constant	Tape thickness, mils					
	0.5	1.0	2.0	4.0		
С	2.270	1.563	1.981	1.959		
g	2.537×10 ⁻¹⁵	3.819×10 ⁻¹⁰	3.240×10 ⁻¹⁵	4.420×10 ⁻¹⁴		
u	-1.000	-0.811	-0.500	-0.345		

Table II.2. - Supermalloy - Uncut

Constant	Tape thickness, mils					
	0.5	1.0	2.0	4.0		
С	1.951	2.000	2.090	2.040		
g	7.48×10 ⁻¹³	4.430×10 ⁻¹³	3.259×10 ⁻¹⁵	1.124×10 ⁻¹⁴		
u	-0.408	-0.825	-0.561	-0.297		

T ble II.3. - Square Permalloy 80 - Uncut

Constant	Tape thickness, mils					
	0.5	1.0	2.0	4.0		
С	0.985	1.713	1.704	1.739		
g	-1.845×10 ⁻⁶	2.759×10 ⁻¹²	1.201×10 ⁻¹²	9.202×10 ⁻¹³		
u	-0.349	-0.225	-0.226	-0.258		

Table II.4. - Ferrite

Constant	Material 3B7
С	2.64
g g	1.544×10 ⁻¹⁵
u	-1.6

TABLE III. - TABULATION OF THE CONSTANT AND EXPONENT

VALUES USED IN EQUATION (4.5)

Table III.1. - Supermalloy - Cut Core

Constant		Tape thickness, mils				
	0.5	1.0	2.0	4.0		
d	1.210	1.307	1.467	1.670		
z	2.038	2.084	1.948	1.948		
s	2.876×10 ⁻¹²	1.065×10 ⁻¹²	7.99×10 ⁻¹³	3.231×10 ⁻¹³		

Table III.2. - Supermalloy - Uncut

Constant	Tape thickness, mils					
	0.5	1.0	2.0	4.0		
đ	1.230	1.300	1.53	1.750		
z	1.940	2.000	2.110	2.020		
8	4.297×10 ⁻¹²	8.676×10 ⁻¹³	9.578×10 ⁻¹⁴	1.056×10 ⁻¹³		

Table III.3. - Square Permalloy 80 - Uncut

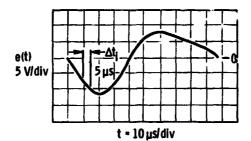
Constant	Tape thickness, mils				
	0.5	1.0	2.0	4.0	
d	0.624	1.456	1.459	1.613	
z	.976	1.696	1.722	1.884	
s	1.718×10 ⁻⁶	7.054×10 ⁻¹²	7.330×10 ⁻¹²	9.447×10 ⁻¹³	

Table III.4. - Ferrite

Constant	Material	
	3B7	
đ	1.04	
z	2.67	
S	3.024×10 ⁻¹³	

The use of Equations (4.4) and (4.5) is not restricted to a computer algorithm. By using a simple hand calculator and the analysis methods described in Section 3.4 the SCL for a particular magnetic material and tape thickness may be easily determined. An example of this process is shown in Figure 20. For the waveform shown the positive half cycle frequency is 8.77 kHz, and the negative half cycle frequency is 12.5 kHz. The measured SCL for this waveshape is 0.144 watts/pound. If the SCL is calculated using 5 µs intervals, the calculated value is 0.132 watt/pound. This represents an error of 8.3 percent. This error will be reduced by decreasing the sample interval.

A computer algorithm was developed which automates the analysis of the large number of sample points encountered in applying this method. A listing of the program is shown in Figure 21. The input/



Waveform excitation into a 1/2 mil Supermalloy cut core

SCL =
$$2.537 \times 10^{-15} (d\phi/dt)^{2.27} f_{eq}^{-1}$$
 (4.3)

	Internal	e(t) - V	døldt - maxwell/sec	SCL watt/fb
	(1	0	0 ,	0
	2	2.5	25x10 ⁶	. 013
-ive half cycle	3	6.5	65x10 ⁰	. 110
	3 4	9.5	95x10 ⁰	. 261
) 5	11.5	115x1v ^o	. 403
	6	11.0	110x10 ⁰	. 364
	6 7 8	8. 5	85x10 ⁰	. 203
	(8	5	5x10 ⁶	0.000
	(1	0	0 ,	0
+ive half cycle	2	0 3	30x10 ⁶	. 027
	$\begin{bmatrix} 1 \\ 2 \\ 3 \end{bmatrix}$	6. 5	65x10 ⁰	. 157
	4	8. 2	82x10 ⁶	. 266
		8.0	80x10 ^c	. 297
	5 6 7	7.0	70x10 ^b	. 186
	ፈ 7	6, 5	65x10 ⁰	. 157
	8	5, 8	58x10 ⁶	. 121
	9	5.0	50x10 ⁶	. 087
	10	3.5	35x10 ⁰	. 039
	ii	3.0	30x10 ⁰	. 027
	12	1.0	10x10 ⁶	. 002
	13	0	0	0

Using equation 3.16:

$$SCL_{T} = \frac{\frac{1.354}{2x8x12.5x10^{3}} + \frac{1.3\%}{2x13x8.77x10^{3}} = 0.132 \text{ watt' pound}}{96.9x10^{-6}} = 0.132 \text{ watt' pound}$$

Figure 20. - An example to demonstrate the application of method 3 and the specific core loss equations.

```
0000108
0000200
0000306 C
0000500 C
0000500 C
0000700 C
                                                                                                                                                               DIMENSION V(2,100), DPDT(2,100), SFL(2,100)
DATA SUPER, SUPERM, SUPPTH, FERT/'S', 'D', 'A', 'F'/
NUMBER OF TUPMS ON TRANSFORMER PRIMARY
READ(5,100) T
NUMBER OF DATA POINTS TOR + PRIMARY VOLTAGE
READ(5,110) J
NUMBER OF DATA POINTS FOR - PRIMARY VOLTAGE
READ(5,110) K
VALUE OF DATA INCREMENT IN MICROSECONDS
READ(5,100) TIME
                                                                                                                                               READ(S,110) J
RUMBER OF DATA POINTS FOR - PRIMAPY VOLTAGE
READ(S,110) K
VALUE OF DATA INCREMENT IN MICROSECONDS
READ(S,100) TINC
TJSJ
TK-K
FOM=1,/(2,-TINC=TJ=1,E-6)
FOFF=1,/(2,-TINC=TJ=1,E-6)
FOFF=1,-(2,-TINC=TJ=1,E-6)
FOFF=1,-(2,-TINC=T
     0000900 C
  0001100
0001200
0001300
0001500
0001500
0001600
     ##91799
0991#G8
0001900
  00018G8
0001900 10
9002900 C
9002100
9002100
9002390
0002500 C
9002500
9002500
  0002700
0002800
       0002900 C
     0003100
  0003200
0003200
0003300
0003500
0003600
0003700
       0003806
       9993900
40
                                                                                                    50
       0005100
                                                                                                    51
          PO05200
  0005200
0005300
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  nen6300
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OF POOR QUALITY
     0007000 70
       0007100
       0007300
  COC7600
0C07500
0C07600
0007700
0007700
0007900
COC200
COC200
COC200
  0002200 100
0002200 100
00023 0 110
00023 0 110
0002400 115
00028700 120
00029700 230
00029700 250
00029700 250
00029700 250
                                                                                                                                                               "RITE(6, 520)TGCL
FORMAT(3)
FORMAT(3)
FORMAT(3)
FORMAT(4)
FORMAT(7)
FORMAT(1)
FORMAT(1
000000 300
000000 370
000000 370
000000 370
000000 370
000000 500
000000 500
000000 500
001000 500
001000 610
001000 620
001000 630
          0016400 630
0010500 640
                                                                                                                                                                            STOP
```

Figure 21. - Program listing for SCL analysis program.

output formats are shown in Figure 22.

The same example shown in Figure 20 was recalculated using the computer algorithm with the sample interval decreased to 2 μs . The calculated SCL for this example was calculated as 0.142 watt/pound. This represents an error of only 1.4 percent when compared with the experimentally measured SCL of 0.144 watt/pound.

This example demonstrates that the third method developed in Section 3.4 may be used to determine the SCL for highly distorted waveforms. Also, by decreasing the sample interval the accuracy of the analysis will be improved.

4.3 Examples

In this section five examples are presented which demonstrate the accuracy of the incremental dolds method. Three magnetic materials were used in the analysis with experimental measurement of their SCL made under both sinusoidal and nonsinusoidal excitation. The waveforms used in these analyses are shown in Figure 23. Both ferrite 3B7 material and Square Permalloy 80 were analyzed using these waveforms. The results of these SCL analyses are shown below:

Material	Sinusoidal excitation		Nonsinusoidal excitation	
	Calculated	Measured	Calculated	Measured
Ferrite	0.388 W/lb	0.369 W/1b	0.244 W/1b	0.247 W/1b
Square Permalloy 80	.777 W/1b	.716 W/1b	.889 W/1b	.831 W/1b

The percentages of error shown for these examples are consistent with thos. measured for the previous examples. Therefore this method may be used for analysis over the wide range of excitation conditions encountered in the application of the magnetic materials.

NUMBER OF PRIMARY TURNS . 0.1000000E 02 POSITIVE VOLTAGE EQUIVALENT FREQUENCY= 0.1000001E 04 HZ +IVE VOLTAGE-VOLTS DP/DT-MAX/SEC 0.1763999E 00 0.1764000E 07 0.5132999E 00 0.5133001E 07 0.7996002E 07 0.7995999E 00 0.1007000E 01 0.1007000E 08 0.1117000E 01 0.1117000E 08 IMPUT FORMAT 0.1117000E 01 0.1117000E 08 f of turns 0.1007000E 01 0.1007000E 08 • 12 of +ive sample PTS 0.7935999E 00 0.5132999E 00 0.7996002E 07 • 12 # of -ive sample PTS 0.5133001E 07 2. 261 Δt in μsec 0.1763999E 00 6.17646COE 07 .765 1.218 NEGATIVE VOLTAGE EQUIVALENT FREQUENCY= 0.1000001E 04 HZ 1.587 # 1.848 # 1.983 +e(t) values -IVE VOLTAGE-VOLTS DP/DT-HAX/SEC # 1.983 (1.848 0.1763909E 00 0.5132999E 00 0.17640005 + 7 G.5133001E 07 1.587 # 1.218 0.7995999E 00 0.7996002E 07 .765 0.1007000E 01 0.1007606E 08 .261 . 0.1117000E 01 0.1117000E 08 .261 0. 1117000E 01 0.1117000E 08 . 765 0.1007000E 01 0.1007000E 03 1.218 0.7995999E 00 0.7996002E 07 . 1.587 0.5132999E 00 0.5133001E 07 # 1.848 0.1763999E 00 0.1764000E 07 1,983 -e(t) values 1,983 # 1.848 1.587 SPECIFIC CORE LOSS (SCL) FOR 0.5E 00 HIL SUPERIALLOY HATERIAL 1.218 .765 material type CONVERTER OPERATING FREQUENCY= 0.1000001E 04 HZ tape thickness h.430e-13 a - constant SCL ON -O.G189252E-OS WATT-SEC/LB-CYCLE .825 SCL OFF-0.6189252E-05 WATT-SEC/LB-CYCLE

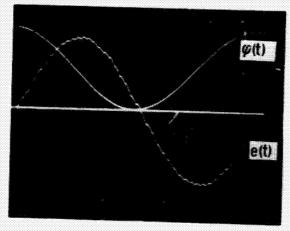
OUTPUT FORMAT

INPUT DATA FOR HAGMETIC WAVEFORM ANALYSIS

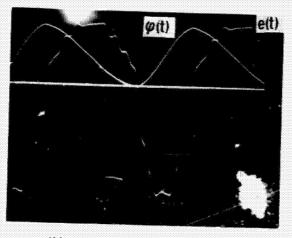
Figure 22. - Input-output formats for automated calculation of SCL of a magnetic material.

0.1237851E-01 WATT/LB

TOTAL SCL-



(a) Sinusoidal excitation.



(b) Nonsinusoidal excitation.

Figure 23, - Waveform excitation for sinusoidal and nonsinusoidal conditions.



This method may also be used with the square wave SCL characteristics to determine the SCL characteristics for sinusoidal excitation. Figure 24 shows the results of an analysis of a 1000 Hz sine wave for various B_{M} levels. The calculated results shown as a bold line are compared with the measured sinusoidal characteristics for 1/2 mil Supermalloy material. The results show that an extremely accurate SCL characteristic may be achieved by this method of analysis.

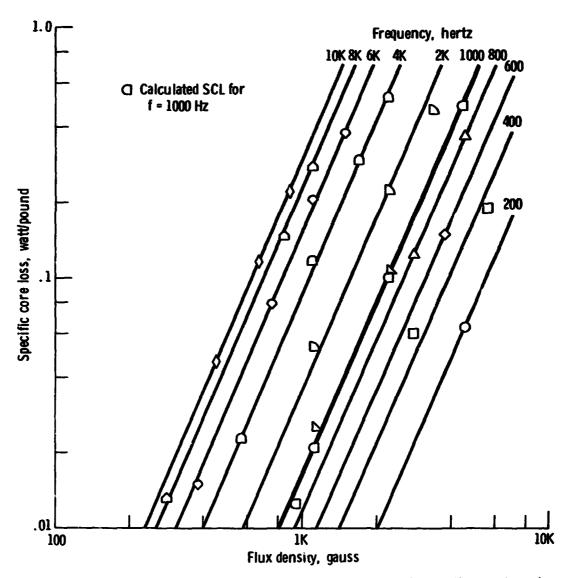


Figure 24. - Comparison of calculated and measured results for specific core loss of 1/2 mil cut Supermalloy under sinusoidal excitation at 1000 Hz.

CHAPTER V

CONCLUSIONS AND SUGGESTIONS FOR FUTURE RESEARCH

5.1 Conclusions

The research covered in this dissertation provides the power electronics designer with a unique analytical tool for investigation of the core loss characteristics of magnetic materials. Analysis methods previously available were based strictly on sinusoidal excitation voltages. Core loss analysis has been based on the assumption that eddy current and hysteresis loss components are separable. By using a novel excitation technique employing a pulse width modulated square wave, a true separation of these two loss components was shown experimentally. This approach provides a straightfurward analysis technique to characterize the basi: core loss process of new materials (Ref. 27).

For fixed B_M conditions the width of the B-H loop was shown to be a direct function of $d\phi/dt$. If this concept is expanded, a family of B-H curves for various B_M levels may be constructed. Using this family, an operating B-H locus may be constructed for any type of waveform excitation. This approach was used to explain the observed difference in B-H width when excited from either a voltage source (narrow B-H width) or a current source (wider B-H width). The flux characteristics under these two source conditions directly influence the observed B-H shape. For sinusoidal voltage excitation the flux is cosinusoidal in shape and has variable $d\phi/dt$ values over one cycle. For sinusoidal current excitation the flux approaches a

square wave. This excitation causes correspondingly higher $d\phi/dt$ values than those encountered for cosinusoidal flux conditions.

These results demonstrate the advantage of using square wave excitation as an analysis tool to determine core loss characteristics. In addition, with square wave excitation the application of Faraday's law becomes straightforward as $d\phi/dt$ is a constant value. The specific core loss characteristics were experimentally measured by holding $d\phi/dt$ constant and varying frequencies to develop a family of core loss characteristics.

In Chapter II the state-of-the-art instrumentation system is discussed. This system was used to obtain large quantities of core loss data quickly and accurately. Through the use of a high power amplifier, and three different feedback conditions, curoperation rent, voltage, and flux sources were easily simulated. This allowed the magnetic characteristics of a large number of source and excitation conditions to be compared and measured easily. With square wave excitation do/dt rates could be calculated dir from excitation voltage levels. A true rms voltmeter was used to measure the excitation voltage level. Accurate core loss was measured for a large number of data points using a wideband wattmeter. This wattmeter has the capability of measuring the power in a 100 kHz square wave to within 2 percent. Instantaneous energy measurements were obtained with a multiplying oscilloscope. A four channel oscilloscope with both time base and X-Y capabilities allowed simultaneous observation of e(t), i(t), p(t), $\phi(t)$, and the B-H loop characteristics of the particular magnetic material under test.

Also in this chapter four different methods for determining stacking factor were covered. Stacking factor error in the published literature is as great as 40 percent. With the measurement methods used in this chapter to determine stacking factor, results consistent to within 5 percent were obtained. Of the four methods used, the oil displacement method yielded the most conservative value for the stacking factor parameter.

Using the SCL characteristics determined under square wave excitation, and the concept of separation of loss variables at constant frequency presented in Chapter III, three analysis techniques were developed which predict the specific core loss of a magnetic material under square wave excitation. The three methods calculated this core loss on the basis of maximum flux density in the core, rate of change of flux within the core, or upon the incremental rate of change of flux in the core.

In Chapter IV the third core loss method was extended to include the analysis of both sinusoidal and nonsinusoidal excitation conditions. To implement the method quickly and reliably, generalized specific core loss equations of the various magnetic materials were derived using experimental data. Once these equations are derived, the specific core losses for various excitation conditions may be readily calculated using a hand calculator. For the examples analyzed in this dissertation a computer algorithm was written to provide a high speed analysis tool which implements the concepts of method three. Analytical core loss predictions compared well with experimentally measured core loss data. These analysis methods are applicable to a broad range of excitation conditions. In addition, the basic SCL charac-

teristics of any desired waveform excitation may be derived.

5.2 Applications in Power Electronics

Most converters in use today excite the magnetic components with some well defined flux waveshape. The emergence of higher power level requirements has led to the development of certain converters using L-C resonant power stages. In these converters the magnetic devices are excited by pseudo-sinusoidal currents. The analysis methods developed in this dissertation are directly applicable to the analysis of the magnetic material behavior under these excitation conditions.

Since all DC-AC/DC converters operate from some steady-state source of raw power, the magnetic exciting voltage as a function of time is a symmetrical or asymmetrical rectangular wave. The exact voltage pulse width is determined by the duty-cycle control mechanism which is incorporated into the converter. In terms of timing implementation, the available duty cycle control schemes are:

Constant on-time T_n , variable off time T_f

Constant T_f , variable T_n

Constant $(T_n + T_f)$, variable individual T_n and T_f

Variable T_n , T_f , and $(T_n + T_f)$

In the examples shown in Chapter III, the duty cycle directly affects the SCL of the magnetic material. Therefore, 'he magnetic losses associated with these converters are not fixed, but become direct functions of duty cycle. The three SCL analysis methods developed allow detailed performance calculations of the magnetic loss experienced in these converter applications.

5.3 Optimization of Tape Thickness

Using the SCL characteristics of a specific magnetic material

the core loss dependence as a function of tape thickness was investigated. Holding $B_{\underline{M}}$ and $d\phi/dt$ constant for the four tape thicknesses measured, it appeared that an optimum tape thickness existed for each of the three magnetic materials tested.

Previous work (Ref. 28) has shown a similar optimization to exist with respect to grain size of the magnetic material. Since the magnetic materials examined in this dissertation were picked from random material batches with uncontrolled grain size, additional controlled experiments using the square wave SCL characteristics as a measure should be performed to examine this observed condition. If a true optimum tape thickness for a given material exists, the traditional practice of using thinner magnetic tapes for higher frequencies would have to be reevaluated.

APPENDIX A

SUMMARY OF SPECIFIC CORE LOSS DATA FOR FERRITE, SUPERMALLOY,

AND SQUARE PERMALLOY 80 CORE MATERIALS

The magnetic core loss characteristics are shown in this appendix with SCL (watts/lo) as the dependent variable, and frequency (Hz) as the dependent variable. Flux density (gauss) and do/dt (maxwells/sec) are shown as parameters. While the units given for CL are not in international units, they are consistent with those used in published data for sinusoidal core loss characteristics. For the tape wound materials tape thicknesses of 1/2, 1, 2, and 4 mils are presented. The ferrite core loss characteristics are for 3B7 material.

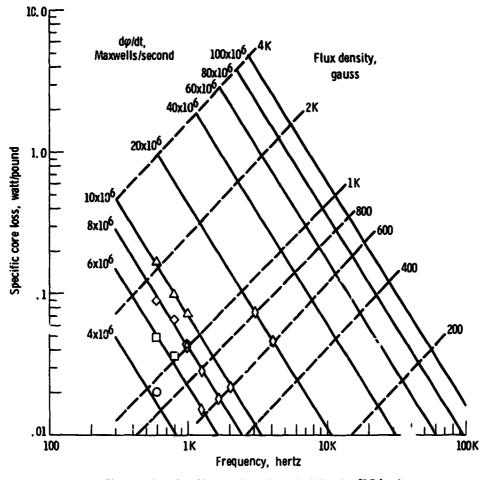


Figure A. 1. - Specific core loss characteristics for 387 ferrite.

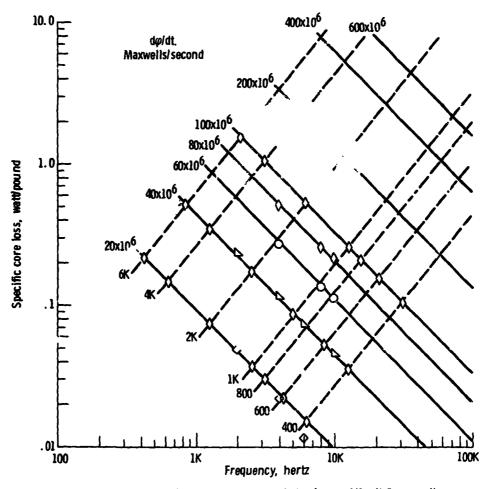


Figure A.2. - Specific core loss characteristics for cut 1/2 mil Supermalloy.

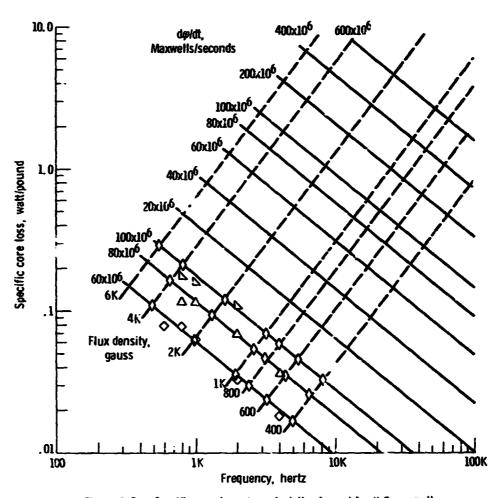


Figure A. 3. - Specific core loss characteristics for cut 1 mil Supermalloy.

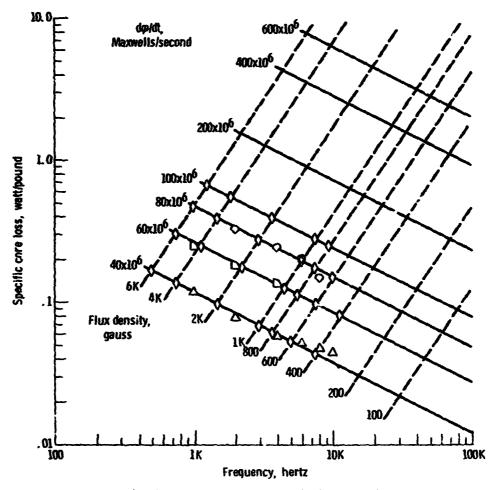


Figure A.4. - Specific core loss characteristics for cut 2 mil Supermalloy.

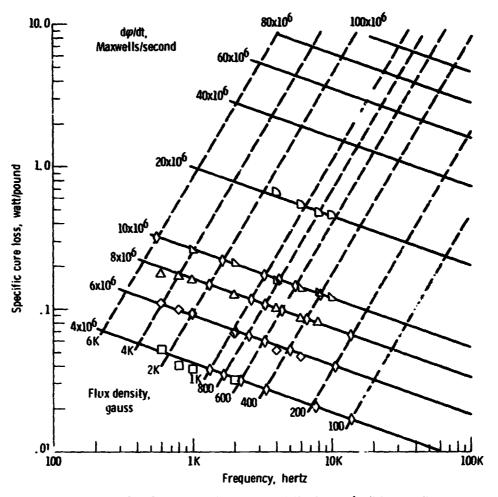


Figure A.5. - Specific core loss characteristics for cut 4 mil Supermalloy.

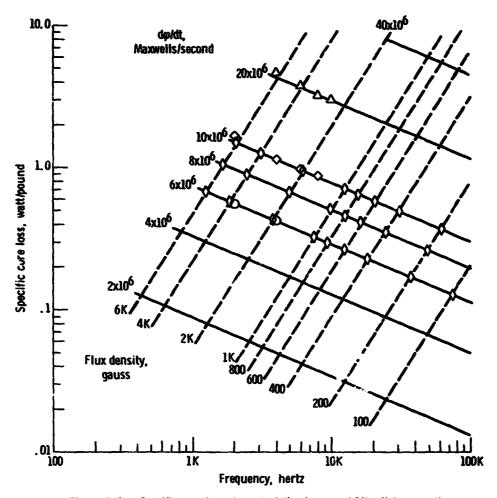


Figure A.6. - Specific core loss characteristics for uncut 1/2 mil Supermatloy.

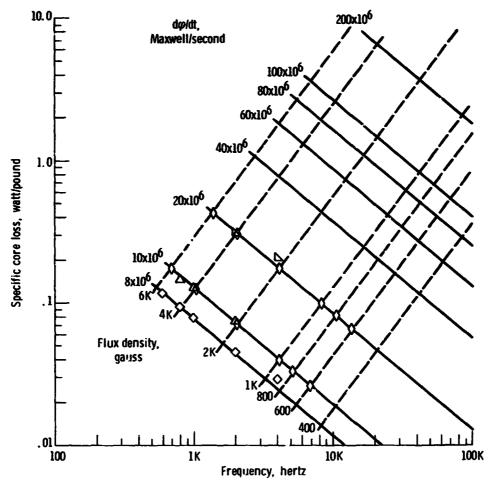


Figure A.7. - Specific core loss characteristics for uncut 1 mil Supermalloy.

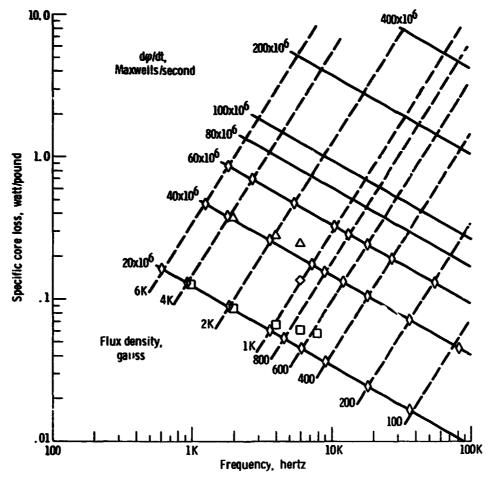


Figure A. 8. - Specific core loss characteristics for uncut 2 mil Supermalloy.

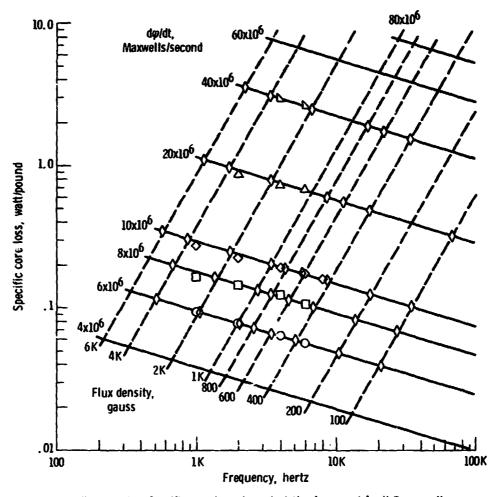


Figure A. 9. - Specific core loss characteristics for uncut 4 mil Supermalloy.

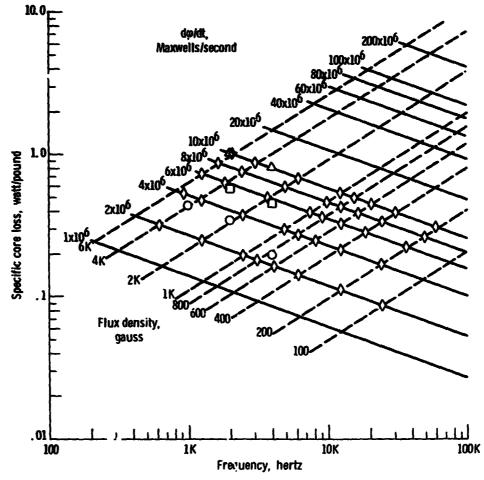


Figure A. 10. - Specific core loss characteristics for uncut 1/2 mil Square Permalloy 80.

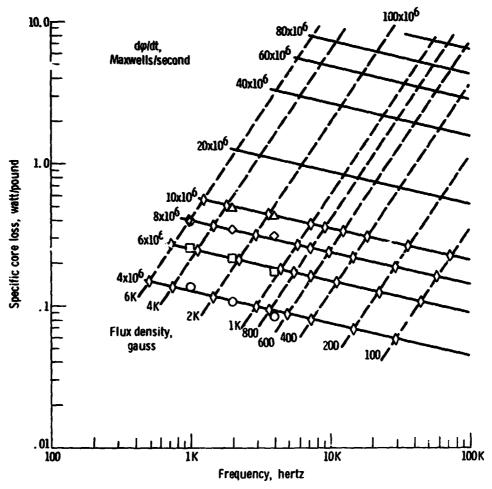


Figure A 11. - Specific core loss characteristics for uncut 1 mil Square Permalloy 80.

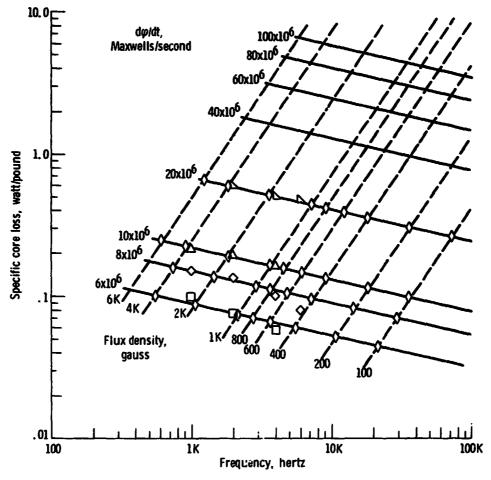


Figure A. 12. - Specific core loss characteristics for 2 mil Square Permalloy 80.

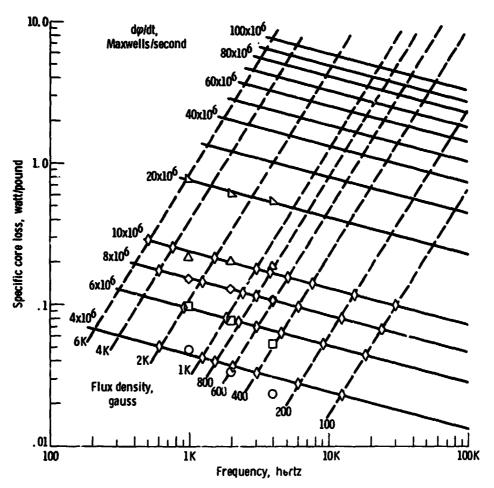


Figure A. 13. - Specific core loss characteristics for 4 mil Square Permalloy 80.

APPENDIX B

GLOSSARY OF SYMBOLS

physical cross sectional area of core, ${\rm cm}^2$ $^{\rm A}_{\rm c}$ effective cross sectional area of core, cm² flux density, gauss maximum flux density, gauss $\mathbf{B}_{\mathbf{M}}$ rate of change of flux, maxwells-sec-1 dø/dt e(t) excitation voltage, V core voltage, V e e_i(t) input voltage to core sample, V incremental energy, W-sec-lb⁻¹-cycle⁻¹ ΔE_{i} total energy, W-sec-1b⁻¹-cycle⁻¹ E f frequency, Hz positive equivalent frequency, Hz negative equivalent frequency, Hz f_{ex} excitation frequency, Hz Н magnetizing force, i_i(t) input current to core, A k transformer equation constant mean magnetic path, cm L negative half cycle increments number of turns Ν positive half cycle increments

input power to core, W

 $p_i(t)$

ρ _{core}	density of core, gm/cm ³
ρ_{oil}	density of oil, gm/cm ³
ф	flux, maxwells
$\phi_{\mathbf{M}}$	maximum flux, maxwells
φ(t)	instantaneous flux, maxwells
SCL	specific core loss, W-lb ⁻¹
scl _i	incremental specific core loss, $W-lb^{-1}$
SCLT	total specific core loss, W-1b ⁻¹
SF	stacking factor
Δt _i	incremental time interval, sec
V core	total volume of core structure, cm ³
v _{Fe}	volume of iron, cm ³
W _{core} air	weight of core in air, gm
W _h	hysteresis loss, J-cycle ⁻¹

displacement weight, gm

ΔW

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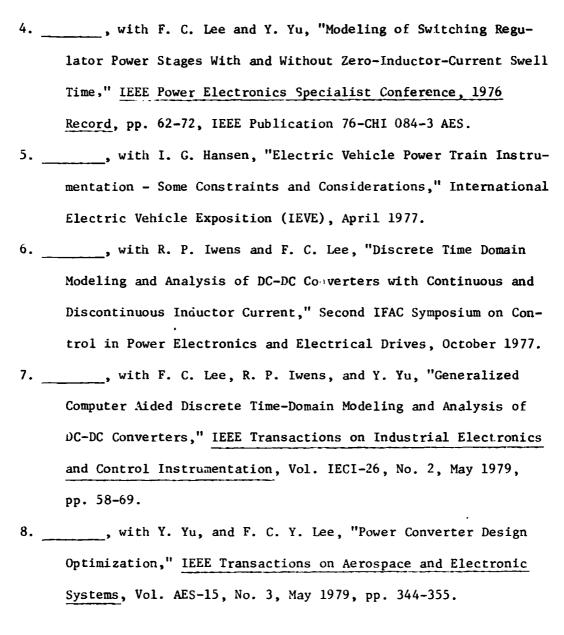
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